Impedance Computation in the Frequency Domain



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Workshop "Ultrafast Beams and Applications" 02-05 July 2019, CANDLE, Armenia



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Motivation



Conventional wakefield computation

$$W_{\parallel}(r,s) = \frac{1}{Q} \int dz \, E_z(r,z,t(s,z))$$

$$t(s,z) = \frac{s+z}{c}$$

- Solve Maxwell's grid equations in the time domain
 - FIT/Cartesian grids/ Dispersion-free methods
 - Co-moving computational window
 - Indirect integration
 - Wall losses by convolution of surface imp.
- Impedance by Fourier transform

$$Z_{\parallel}(r,\omega) = -\frac{1}{c} \frac{1}{\tilde{\lambda}(\omega)} \int ds W_{\parallel}(r,s) \exp\left(-i\frac{\omega}{c}s\right)$$



Motivation



- Long range wakefields
 - Low frequency, long bunches, bunch trains, wall heating
- Approximation of geometry
 - Geometrical details smaller than bunch length, smooth tapering etc.
- Dispersive problems
 - Surface impedance, dielectrics
 - Free-space and waveguide boundary conditions
- Radiation effects
 - Curved beam trajectories (and CSR)
 - Accelerated beams
- Coupler and Pickup signals





The frequency domain problem

$$\nabla \times \mu^{-1} \nabla \times E - k_0^2 \varepsilon E = -jk_0 Z_0 J_s \qquad J_s(x, y, z, \omega) = \rho(x, y) e^{-i\frac{\omega}{\nu} z}$$

• Weak FEM formulation: find $E \in H(curl)$ such that:

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h = -jk_0 Z_0 \int dV J_s \cdot v_h$$

$$+ \oint_s dS \ n \cdot [v_h \times \mu^{-1} \nabla \times E]$$
boundary term
$$\forall v_h \in H(curl)$$





Treatment of boundary surfaces

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h =$$

$$-jk_0 Z_0 \int dV J_s \cdot v_h + \int_{S_{SIBC}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E] + \int_{S_{WG}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E]$$
resistive wall in & outgoing pipes

Resistive wall boundary

$$\oint_{S_{SIBC}} dS \ n \cdot [v_h \times \mu^{-1} \nabla \times E] = \dots = j \omega \mathbf{Y}_{\mathcal{S}}(\boldsymbol{\omega}) \oint_{S_{SIBC}} dS \ v_h \cdot [n \times n \times E]$$

Simple modification of the system matrix on SIBC surfaces No fitting of the surface impedance function or ADE/convolution is needed





Treatment of boundary surfaces

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h =$$

$$-jk_0 Z_0 \int dV J_s \cdot v_h + \int_{S_{SIBC}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E] + \int_{S_{WG}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E]$$
resistive wall in & outgoing pipes

Beam pipe boundaries

$$\begin{split} n \times \nabla \times E &= n \times \nabla \times E^{inc} + \sum_{m} a_{m}^{TE} \gamma_{m}^{TE} e_{m}^{TE} + \sum_{m} a_{m}^{TM} \frac{-k_{0}^{2}}{\gamma_{m}^{TM}} e_{m}^{TM} \\ a_{m}^{TE} &= \int_{S_{WG}} dS e_{m}^{TE} \cdot \left[E - E^{inc} \right] \\ \text{Reflection coefficients for each mode} \\ a_{m}^{TM} &= \int_{S_{WG}} dS e_{m}^{TM} \cdot \left[E - E^{inc} \right] \end{split}$$





Beam pipe boundary conditions

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon \ E \cdot v_h + \sum_m P_m^{TE}(E) + \sum_m P_m^{TM}(E) = -jk_0 Z_0 \int dV J_s \cdot v_h + \oint_{S_{WG}} dS \ n \cdot [v_h \times \mu^{-1} \nabla \times E^{inc}] + \sum_m U_m^{TE} + \sum_m U_m^{TM}$$

with $P_m^{TE}(E) = -\gamma_m^{TE} \left(\int_{S_{WG}} dS \ v_h \cdot e_m^{TE} \right) \left(\int_{S_{WG}} dS \ e_m^{TE} \cdot E \right), \qquad P_m^{TM}(E) = \cdots$

and matrix representation (TE):

$$P_m^{TE}(E) \rightarrow P_m^{TE} \cdot \mathbf{e} = -\gamma_m^{TE} \mathbf{R}^T \cdot \mathbf{M}_m^{TE} \cdot \mathbf{R} \cdot \mathbf{e} \qquad [\mathbf{R}]_{ij} = \int_{S_{WG}} dS \, \varphi_i^{2D} \cdot \varphi_j^{3D}$$

$$\mathbf{M}_m^{TE} = \mathbf{e} \, \frac{TE}{m} \otimes \mathbf{e} \, \frac{TE}{m} \qquad \text{dense modal dyadic} \qquad 3D\text{-to-2D projection matrix}$$





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- Beam pipe boundary excitation
- For an ultra-relativistic bunch (same idea for $\beta < 1$):



2D-electrostatic problem at both ends of the pipe

- Modal contribution to the RHS

$$U_m^{TE}(E^{inc}) = -\gamma_m^{TE} \left(\int_{S_{WG}} dS \, v_h \cdot e_m^{TE} \right) \left(\int_{S_{WG}} dS \, e_m^{TE} \cdot E^{inc} \right)$$

$$U_m^{TE}(E^{inc}) \rightarrow U_m^{TE} \cdot \mathbf{e}^{inc} = -\gamma_0^{TE} \mathbf{R}^T \cdot \mathbf{M}_m^{TE} \cdot \mathbf{R}^{2D} \cdot \mathbf{e}^{inc}$$

...do this for all waveguide modes supported in the pipe



Collimator – hybrid meshes 80 Re (E_z) - tetrahedral mesh 60 Re (E_z) - mixed mesh TET HEX Im (E_z) - mixed mesh 40 **PYRA** 20 E_z / a.u. 0 -20 -40 Longitudinal -60 wakefield on axis -80 at 100MHz -100 -1.5 -1.2 -0.9 -0.3 0 0.3 0.6 0.9 1.2 -0.6 1.5 z/m

Frequency Domain Formulation

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Resistive Wall Impedance







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 Equivalent circuit representation in/outgoing modal waves b_1 a_1 Z_2 b_2 a_2 Z_1 l_2 l_b u_2 (N+1)-port l_1 network u_1 u_b $\rightarrow b_N$ a_N ω_0^1 a_b

- Only synchronous field prop. in the +z direction contributes to beam voltage
 - Beam as a one-way transmission line of infinite length



 $b_{b} = 0$

 $u_b + Z_b i_b$

 $a_b =$



S-Matrix / hybrid matrix representation

$$\begin{pmatrix} S & k \\ h & r \end{pmatrix} \begin{pmatrix} a_m \\ a_b \end{pmatrix} = \begin{pmatrix} b_m \\ 0 \end{pmatrix} \quad \Rightarrow \quad \begin{pmatrix} S & \tilde{k} \\ \tilde{h} & Z_b \end{pmatrix} \begin{pmatrix} a_m \\ i_b \end{pmatrix} = \begin{pmatrix} b_m \\ -u_b \end{pmatrix}$$

- Concatenation for cascaded structures:

Matching conditions:



- Reduces to total scattering matrix:

$$\begin{pmatrix} S^{tot} & \tilde{k} \\ \tilde{h} & \mathbf{Z}_{b}^{tot} \end{pmatrix} \begin{pmatrix} a_{m} \\ i_{b} \end{pmatrix} = \begin{pmatrix} b_{m} \\ -u_{b}^{tot} \end{pmatrix}$$





Two cavity / two mode coupling example







Tesla 1.3GHz cavity







Tesla 1.3GHz cavity









Tesla 1.3GHz cavity







Dechirper (LCLS, E-XFEL)



Single period computation layout





Analytical (Bane et al.)

03.07.2019 | TU Darmstadt | Fachbereich 18 | Institut Theorie Elektromagnetischer Felder | PD Dr.rer.nat. Erion Gjonaj | 19



2.0

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Impedance vs. no. periods

Impedance vs. coupling modes



Curved beam trajectories



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E-XFEL BC0





Curved beam trajectories







Summary & Conclusions



- The frequency domain approach
 - Fills the gap for some important wakefield/impedance problems
 - Complicated chamber geometry, long range / low frequency fields, resistive, rough surfaces, dispersive materials, beam signals on waveguide openings

- FEM Frequency domain formulation

- Beam port boundary conditions
- Mixed mesh discretization
- Concatenation using generalized S-Matrix formulation
 - Efficient /accurate impedance computation of large cavity chains
 - Periodic structures (dechirper)
- Impedance (+wall losses/radiation) computation for curved beam trajectories
- Limitation: huge size of discrete problem for ultra-high frequencies
 - Domain decomposition
 - Parallel multigrid solvers
 - Fast frequency sweeps and spectral evaluation by model order reduction,...



Thank You for your attention