# Synchronous acceleration with tapered dielectric-lined waveguides

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# Acceleration of non-relativistic particles in high frequency structures What is the challenge?

Energy gain per rf period for a constant phase

$$\Delta \gamma = \int_{0}^{kz=2\pi} \frac{eE(z,t)}{mc^{2}} dz = \int_{0}^{kz=2\pi} \frac{e}{mc^{2}} E_{0} \left( \sin(\phi) + \frac{\sin(\phi+2kz)}{\gamma} \right) dz$$

only in standing wave structures, does not contribute

$$\Delta \gamma = 2\pi \frac{eE_0}{mc^2 k} \sin(\phi) = 2\pi\alpha \sin(\phi)$$

The normalized vector potential  $\alpha = \frac{eE_0}{mc^2k}$  describes the longitudinal beam dynamics independent of the rf wavelength.

For  $\alpha > 1$  a particles starting at rest reaches relativistic energies within one rf period.

For  $\alpha < 1$  phase slippage is significant. Efficient acceleration requires matching of the phase velocity of the wave to the (increasing) velocity of the particle.

# The normalized vector potential shrinks with increasing frequency

... because the gradient is not increased as much as the wavelength decrease



Th. Vinatier et al. EAAC 2017, NIM A 909 (2018) 185–192

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# Synchronous acceleration in a tapered dielectric waveguide



Generalized TM<sub>01</sub> mode description:

$$E_{z} = E_{0}I_{0}(rk_{r})\sin\left(\omega t - \int_{0}^{z}k_{z}dz + \varphi\right)$$

 $E_0$ ,  $k_r$  and  $k_z$  are functions of the longitudinal coordinate

# Synchronous acceleration in a tapered dielectric waveguide



average particle energy

IT WORKS:

The electron stays on a constant phase and thus gains continuously energy.

BUT:

The fields are nonlinear; not due to the tapering but due to the matching to velocities below speed of light:

THIS IS FUNDAMENTAL!

Example: Energy gain in a tapered dielectric structure

# Synchronous acceleration in a tapered dielectric waveguide

z = 9.9900E - 02 m



Particles at large radii see higher fields...

First order approximations:

$$F_z = eE_0\sin(\varphi)\left(1 + \frac{k_r^2 r^2}{4}\right)$$

$$F_r = eE_0\cos(\varphi)\frac{k_0}{\gamma^2\beta}\left(\frac{r}{2} + \frac{k_r^2r^3}{16}\right)$$

- the beam size has to be small in comparison to the wavelength
- higher initial energies are beneficial

# Stable operation when slightly over-powered...



- structure is matched for a gradient of 100 MV/m
- initial energy 100 keV, 10 cm DLW
- broad phase acceptance when the gradient is somewhat higher than the design gradient

Energy gain vs. start phase for different operation gradients

# .... strong (on-axis) bunch compression when slightly overpowered



End phase vs. start phase for different operation gradients

# **Schematic layout of compact THz accelerator**



# **Compact RF Gun**

design and construction by group of Franz Kaertner



- 3 GHz single cell cavity with cathode tip and choke mode filter
- > driven by a compact solid state amplifier (10 kW)
- high field (100 MV/m) on the cathode but low end energy (150 – 200 keV)

# **Prototype of the Compact RF Gun**



# **Simulation results for the Compact Gun**



- > 150 fC bunch
- ➤ 200 keV energy gain
- ➢ Negative phase injection → bunch compression
- Transverse emittance:0.1 mm mrad

<sup>...25</sup> µm bunch length....

# Simulation results Gun & DLW & Focusing

initial energy 200 keV, frequency of the DLW 300 GHz, matched gradient 100 MV/m,  $\alpha = 0.03$ 



Inverse bunch length (top) and transverse emittance (bottom) as function of gradient and start phase.

White contour lines show the charge transmission.

Parameters for shortest bunch length:

charge: 80 fC

bunch length: 730 nm

emittance: 158 nm

energy spread: 83 keV

# **Bunching mechanism in standard and tapered structures**



- negative starting phase: tail gains more energy than head the bunch is thus compressed; for positive starting phases the bunch is decompressed
- > particles starting below the  $\eta = 1$ line move toward higher phase (from left to right); particles starting above the  $\eta = 1$  move toward lower phase (from right to left)
- fix points I and II are not stable, while fix point III is stable

### **Bunching mechanism in standard and tapered structures**



# **Beam dynamics near fix point III**

A first order Taylor expansion of the phase near the fix point yields the differential equation

$$\Delta \varphi'' + \frac{3\gamma'}{\gamma} \Delta \varphi' + \frac{k_0 \eta \gamma' |\sin \varphi_0|}{\gamma^3} \Delta \varphi = 0$$

(simplified form for  $\beta = 1$ ), with the solution

$$\Delta \varphi(\gamma) = \Delta \varphi_0 \frac{\gamma_0}{\kappa_0 \gamma} \left[ -C_1 J_2(\kappa) + C_2 Y(\kappa) \right]$$

 $C_1$ ,  $C_2$  = constants

 $J_2$ ,  $Y_2$  = Bessel functions

 $\Delta \varphi_0$  = initial phase relative to fix point phase

 $\gamma_0$  = initial energy

$$\kappa = 2\sqrt{\frac{\sqrt{\eta^2 - 1}}{\alpha \gamma}}$$

The relevant parameters to describe the beam dynamics are: the initial energy  $\gamma_0$ , the normalized vector potential  $\alpha$  and the ratio of gradient to matched gradient  $\eta$ .

#### **Proof of relevant parameters**



#### Phase offset vs. energy for different vector potentials



The phase swings over for too low alpha ( $\alpha \le 0.03$ ), i.e. a bunch is compressed, decompressed and compressed again. For large alpha ( $\alpha = 0.27$ ), the damping due to the increased gamma is too strong, the compression is stopped too early  $\alpha = 0.13$  is close to the aperiodic limit.

 $\eta = \text{const}=1.1$ 

# Simulation results Gun & DLW & Focusing

initial energy 200 keV, frequency of the DLW 75 GHz, matched gradient 100 MV/m



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75 GHZ, 100 MV/m: α = 0.12

Parameters for shortest bunch length:

charge: 100 fC

bunch length: 473 nm

emittance: 227 nm

energy spread: 59 keV

Page 19

# **Bunch length during the acceleration process**



# **Final bunch length**

